

# OSCILLATORY FLOW OF CASSON FLUID BETWEEN PARALLEL PLATES WITH AN INCLINED MAGNETIC FIELD

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## Abstract

The study of heat and mass transfer of oscillatory casson flow in porous medium subject to an inclined magnetic field, radiative heat flux and heat source is presented. It is supposed that Casson fluid is little conductive and produced emf is insignificant. The solutions of coupled partial differential equations of velocity, temperature and concentration profiles are found using Galerkins technique of finite element method. The effect of various parameters such as Reynolds number  $Re$ , Grashoff number  $Gr$ , Solute Grashoff number  $Gc$ , Peclet number  $Pe$ , Hartman number  $Ha$ , Schmidt number  $Sc$ , Permeability parameter  $K$ , Radiative parameter  $R$ , Heat generation parameter  $S$ , Chemical reaction parameter  $Kr$  and frequency parameter  $\omega$  on velocity, temperature and concentration are shown graphically and skin friction, Nusselts number and Sherwood number are discussed by tables

**Keywords:** Free convective, Skin-friction, Nusselt number, Sherwood number, Casson fluid.

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## 1 Introduction

Fluid flow with heat and mass transfer in porous medium has many applications on geophysics, petro chemical engineering, oceanography, boundary layer control in aerodynamics and polymer technology. Magneto hydrodynamic flows of electrically conductive fluid plays important role in MHD bearings, power generators, metrology, astrophysics to study the stellar and solar structures, inter stellar matters and radio wave propagation through ionosphere, extraction, electromagnetic propulsion and has industrial applications in molten iron flow, extraction of crude oil, electrostatic precipitation. [1] presented heat transfer in MHD flow between parallel finitely conductive plates in presence of a transverse magnetic field.[2] considered the steady flow of a viscous incompressible fluid bounded by two infinite insulated horizontal plates and heat transfer through it examined.[4] probed the boundary layer memory flow and given analytic solution for heat and mass transfer considering hall current effects.[5] discussed a MHD flow of a viscous fluid between two parallel porous plates. The effect of magnetic field and pressure gradient on velocity and temperature focused.[6]analyzed the MHD flow through parallel plates with heat and mass transfer and showed that Hartman number reduces the flow.[7] ascertained that Hatmann number enhances induced magnetic field and magnetic Reynolds number influenced velocity in their study of heat and mass transfer between vertical parallel plate.[8] investigated MHD Poiseuille flow between two parallel plates with heat transfer an showed that Hartmann number suppresses flow and Prandtle number subsides temperature proffles.[9] discussed the steady Poiseuille flow between two infinite parallel plates in an inclined magnetic field and ascertained magnetic field reduces the velocity.[11] obtained the analytic solutions for velocity,temperature,concentration, skin-friction, heat and mass transfer using Laplace technique.[12] studied MHD free convective non-newtonian flow with variable permeability. Numerical solution of unsteady two dimensional hydromagnetic flow with heat and mass transfer of casson fluid discussed by [13].[14] given observation of unsteady MHD poiseuille flow between two infinite parallel plates with heat and mass transfer subjected to porous material and an inclined magnetic field.[15]discussed heat and mass transport on MHD free convective flow through porous medium past an infinite vertical plate.[16]shown the effect of an inclined magnetic field is to reduce flow of dusty fluid in unsteady case.[17] discussed effects of variable thermal conductivity and heat source/sink on MHD flow near a stagnation point on a boundary conditions.[18] highlighted the heat and mass transfer in steady flow in presence of an inclined magnetic field and found velocity first decreases and then rises gradually.

The aim of the paper is to extend the work of [10] to memory fluid with mass transfer

## 2 Formulation of the problem

Consider the oscillatory casson fluid flow of an electrically conducting fluid and heat transfer between infinite parallel plates at a distance  $d$  apart filled with a porous medium under the influence of an inclined magnetic field and heat source. The  $x^*$ -axis is taken along the plate when ( $y^* = 0$ ) and the  $y^*$ -axis is taken normal to the plate. As the plates are infinite in length, the velocity and temperature fields are functions of  $y^*$  and  $t^*$  only.

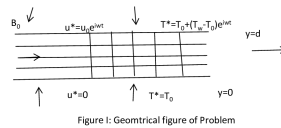


Figure 1: Geometrical figure of Problem

Figure 1: Geometrical representation of problem

Under the above assumption and usual Boussineq approximation, the governing equations of motion, energy and concentration are as follows

$$\frac{\partial u^*}{\partial t^*} = -\frac{1}{\rho} \frac{\partial p^*}{\partial x^*} + \nu \beta^* \frac{\partial^2 u^*}{\partial y^{*2}} + g\beta(T^* - T_0) + g\beta^{**}(C^* - C_0) - \frac{\sigma_e B_0^2 \sin^2 \phi}{\rho} u^* - \frac{\nu}{K^*} u^* \quad (2.1)$$

$$\frac{\partial T^*}{\partial t^*} = \frac{\kappa}{\rho C_p} \frac{\partial^2 T^*}{\partial y^{*2}} - \frac{1}{\rho C_p} \frac{\partial q_r}{\partial y^*} + \frac{Q_0}{\rho C_p} (T^* - T_0) \quad (2.2)$$

$$\frac{\partial C^*}{\partial t^*} = D \frac{\partial^2 C^*}{\partial y^{*2}} - K_r (C^* - C_0) \quad (2.3)$$

The boundary conditions are given by:

$$y^* = 0 : u^* = 0, T^* = T_0, C^* = C_0 \quad (2.4)$$

$$y^* = d : u^* = U_0 e^{i\omega t^*}, T^* = T_0 + (T_w - T_0) e^{i\omega t^*}, C^* = C_0 + (C_w - C_0) e^{i\omega t^*} \quad (2.5)$$

It is assumed that the fluid is optically thin with a relative low density and radiative heat flux is according to [3] and given by

$$\frac{\partial q_r}{\partial y^*} = 4\alpha^2(T^* - T_0) \quad (2.6)$$

Introducing the following dimensionless quantities,

$$u = \frac{u^*}{U_0}, x = \frac{x^*}{d}, y = \frac{y^*}{d}, t = \frac{U_0}{d}t^*, \omega = \frac{d}{U_0}\omega^*, P = \frac{d}{\nu\rho U_0}P^*, \theta = \frac{T^* - T_0}{T_w - T_0}, Re = \frac{U_0 d}{\nu}$$

$$Ha^2 = \frac{\sigma_e B_0^2 d^2}{\nu\rho}, Gr = \frac{g\beta d^2(T_w - T_0)}{\nu U_0}, K^2 = \frac{d^2}{K^*}, Pe = \frac{U_0 d \rho C_p}{\kappa}, R^2 = \frac{4\alpha^2 d^2}{\nu\rho C_p},$$

$$S = \frac{Q_0 d^2}{\kappa}, Sc = \frac{D}{\nu}, K_r^* = \frac{k_r d}{u_0}, Gc = \frac{g\beta_1 d^2(C_w - C_0)}{\nu u_0}, \phi = \frac{C - C_0}{C_w - C_0},$$

into eq (2.1),(2.2) and (2.3) we get

$$Re \frac{\partial u}{\partial t} = -\frac{\partial p}{\partial x} + \beta^* \frac{\partial^2 u}{\partial y^2} + Gr\theta - (M^2 + K^2)u + Gc\phi \quad (2.7)$$

$$Pe \frac{\partial \theta}{\partial t} = \frac{\partial^2 \theta}{\partial y^2} - R^2\theta + S\theta \quad (2.8)$$

$$\frac{\partial \phi}{\partial t} = \frac{1}{Sc} \frac{\partial^2 \phi}{\partial y^2} - K_r\phi \quad (2.9)$$

where  $M = Ha \sin \psi$

The boundary conditions are reduced to

$$y = 0 : u = 0, \theta = 0, \phi = 0 \quad (2.10)$$

$$y = 1 : u = e^{i\omega t}, \theta = e^{i\omega t}, \phi = e^{i\omega t} \quad (2.11)$$

### 3 Method of Solution

For a purely oscillatory flow, substituting

$$-\frac{\partial P}{\partial x} = \lambda e^{i\omega t}, u(y, t) = u_0(y)e^{i\omega t}, \theta(y, t) = \theta_0(y)e^{i\omega t} \text{ and } \phi(y, t) = \phi_0(y)e^{i\omega t} \quad (3.1)$$

into equations (2.7),(2.8) and (2.9) we get

$$u_\theta^{11} - v_1 u_0 = -\frac{1}{\beta^*}(\lambda + Gr\theta_0 + Gc\phi_0) \quad (3.2)$$

$$\theta_0^{11} - \nu_2 \theta_0 = 0 \quad (3.3)$$

$$\phi_0^{11} - v_3\phi_0 = 0 \quad (3.4)$$

Now, the corresponding boundary conditions are

$$y = 0 : u_0 = 0, \theta_0 = 0, \phi_0 = 0 \quad (3.5)$$

$$y = 1 : u_0 = 1, \theta_0 = 1, \phi_0 = 1 \quad (3.6)$$

Equations (3.2), (3.3) and (3.4) are ordinary second order coupled differential equations and are solved under the boundary conditions (3.5) and (3.6). Through straight forward calculations  $u_0$  and  $\theta_0, \phi_0$  are found. Finally, the expressions of  $u(y,t), \theta(y,t)$  and  $\phi(y,t)$  are known as given below

$$u_0 = l_5(y^2 - y) + y + il_6(y^2 - y) \quad (3.7)$$

$$u(y, t) = u_0 e^{i\omega t} \quad (3.8)$$

$$\theta_0 = (m_1 + 1) \cos \omega t - m_2 \sin \omega t \quad (3.9)$$

$$\theta(y, t) = \theta_0 e^{i\omega t} \quad (3.10)$$

$$\phi_0 = [n_1 y^2 + (1 - n_1)y] \cos \omega t - n_2(y^2 - y) \sin \omega t \quad (3.11)$$

$$\phi(y, t) = \phi_0 e^{i\omega t} \quad (3.12)$$

The dimensionless stress tensor in terms of the skin-friction coefficient at upper the plates are given by

$$C_f = (l_5 + 1) \cos \omega t - l_6 \sin \omega t \quad (3.13)$$

The non dimensional rate of heat transfer in terms of the Nusselts number at the upper plate is given by

$$Nu = (m_1 + 1) \cos \omega t - m_2 \sin \omega t \quad (3.14)$$

The non dimensional rate of mass transfer in terms of Sherwood number at the upper plate is given by

$$Sh = (n_1 + 1) \cos \omega t - n_2 \sin \omega t \quad (3.15)$$

Constants  $m_1, m_2, n_1, n_2, l_1 \dots$  are not presented for the sake of brevity.

## 4 Results and Conclusion

The effect of an inclined magnetic field, radiation and heat source on an oscillatory flow of an incompressible Casson fluid through a porous medium between parallel plates are investigated. Equations of momentum, energy and concentration which governs the fluid flow, heat transfer and mass transfer are solved. The effects of various physical parameters on the fluid velocity, temperature and concentration are depicted graphically and skin friction, Nusselt number and Sherwood number at the upper plate are discussed numerically. Table 1 highlights that, Skin friction  $C_f$  dominates with rise in Hartmann number Ha, phase angle  $\psi$ , permeability parameter K, frequency  $\omega$ , radiation parameter R, Schmidt number Sc and chemical reaction parameter Kr. And also, skin friction reduces with the rise in Grashoff number Gr, solute Grashoff number Gc, Reynolds number Re, Peclet number Pea and heat source parameter S.

$\alpha$	Ha	$\lambda$	Re	K	R	S	$\omega$	Kr	Gr	Gc	Pe	Sc	t	$\psi$	Cf
0.25	2	0.1	2	1	2	5	1	1	10	1	0.71	2	0	$\frac{\pi}{4}$	-1.440590
0.25	4	0.1	2	1	2	5	1	1	10	1	0.71	2	0	$\frac{\pi}{4}$	0.540196
0.25	2	0.2	2	1	2	5	1	1	10	1	0.71	2	0	$\frac{\pi}{4}$	-1.44414
0.25	2	0.1	4	1	2	5	1	1	10	1	0.71	2	0	$\frac{\pi}{4}$	-0.382325
0.25	2	0.1	2	2	2	5	1	1	10	1	0.71	2	0	$\frac{\pi}{4}$	-0.425023
0.25	2	0.1	2	1	4	5	1	1	10	1	0.71	2	0	$\frac{\pi}{4}$	0.460438
0.25	2	0.1	2	1	2	6	1	1	10	1	0.71	2	0	$\frac{\pi}{4}$	-1.857449
0.25	2	0.1	2	1	2	5	2	1	10	1	0.71	2	0	$\frac{\pi}{4}$	-0.436621
0.25	2	0.1	2	1	2	5	1	2	10	1	0.71	2	0	$\frac{\pi}{4}$	-1.405197
0.25	2	0.1	2	1	2	5	1	1	15	1	0.71	2	0	$\frac{\pi}{4}$	-3.096498
0.25	2	0.1	2	1	2	5	1	1	10	2	0.71	2	0	$\frac{\pi}{4}$	-1.687013
0.25	2	0.1	2	1	2	5	1	1	10	1	1.0	2	0	$\frac{\pi}{4}$	-1.446566
0.25	2	0.1	2	1	2	5	1	1	10	1	0.71	4	0	$\frac{\pi}{4}$	-1.404351
0.25	2	0.1	2	1	2	5	1	1	10	1	0.71	2	1	$\frac{\pi}{4}$	0.543622
0.25	2	0.1	2	1	2	5	1	1	10	1	0.71	2	0	$\frac{\pi}{3}$	-0.954435
0.5	2	0.1	2	1	2	5	1	1	10	1	0.71	2	0	$\frac{\pi}{4}$	-1.440590

Table 1: Numerical Values of skin friction coefficient at the upper plate for the various values of physical properties

Table 2 shows ,the rate of heat transfer i:e Nusselt number Nu increases with increase in radiation parameter R and frequency  $\omega$  .And decreases with increase in heat source parameter S and time t.

R	S	Pe	$\omega$	t	Nu
2	5	0.71	1	0	0.582642
4	5	0.71	1	0	1.113269
2	6	0.71	1	0	0.423934
2	5	1.0	1	0	0.588231
2	5	0.71	2	0	0.599597
2	5	0.71	1	1	0.226686

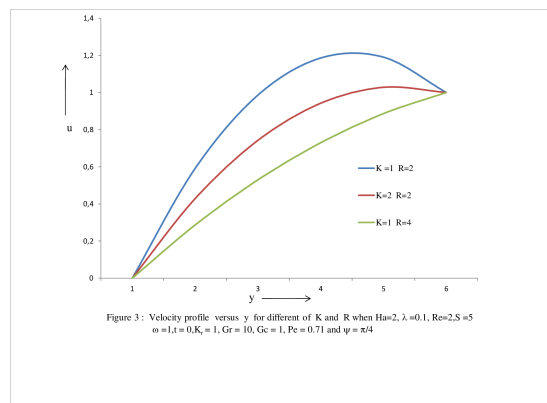
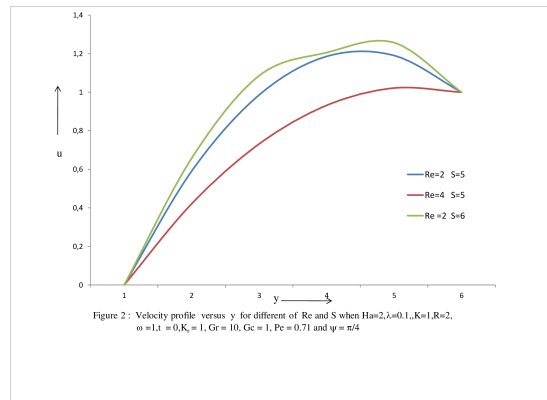
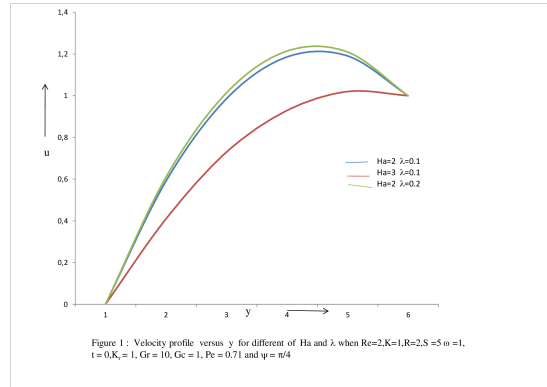
Table 2: Numerical Values of rate of heat transfer i:e Nusselt number at the upper plate for the various values of physical properties

From Table 3 we can conclude that ,the Schimdt number Sc,chemical reaction parameter Kr,frequency  $\omega$  influences the mass transfer .And also,with time t mass transfer dips.

Sc	Kr	$\omega$	t	Sh
2	1	1	0	1.43189
4	1	1	0	1.707547
2	2	1	0	1.72500
2	1	2	0	1.50000
2	1	1	1	0.66385

Table 3: Numerical Values of rate of mass transfer i:e Sherwood number at the upper plate for the various values of physical properties

From Figures 1-7 it is observed that ,velocity shoots up for higher values of Grashoff number Gr,Solute Grashoff number Gc,heat generation parameter S .whereas,velocity has reverse effect for low values of Hartmann number Ha,Reynolds number Re,Radiation parameter R,frequency  $\omega$ ,permeability parameter K. Peclet number Pe and chemical reaction parameter Kr has no affects on velocity.



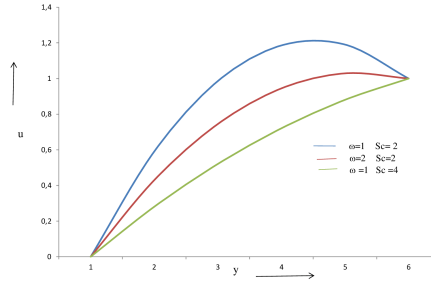


Figure 4: Velocity profile versus  $y$  for different of  $\omega$  and  $Sc$  when  $Ha=2, \lambda=0.1, Re=2, S=5, R=2, \tau=0, K=1, Gr=10, Gc=1, Pe=0.71$  and  $\psi=\pi/4$

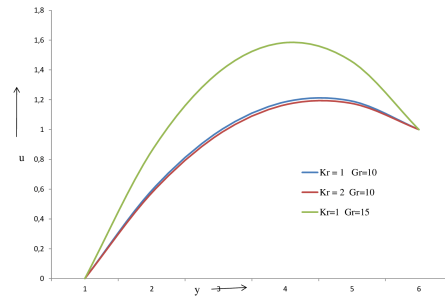


Figure 5: Velocity profile versus  $y$  for different of  $Kr$  and  $Gr$  when  $Ha=2, \lambda=0.1, Re=2, S=5, R=2, \tau=0, \omega=1, Sc=2, Gc=1, Pe=0.71$  and  $\psi=\pi/4$

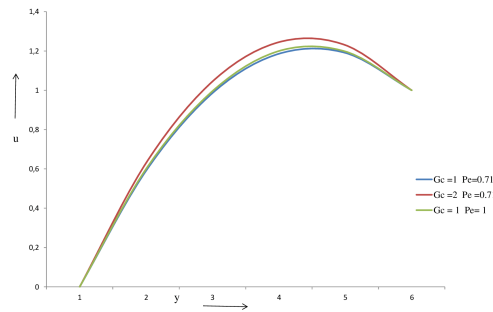


Figure 6: Velocity profile versus  $y$  for different of  $Gc$  and  $Pe$  when  $Ha=2, \lambda=0.1, Re=2, S=5, R=2, \tau=0, K=1, Gr=10, \omega=1, Sc=2$  and  $\psi=\pi/4$

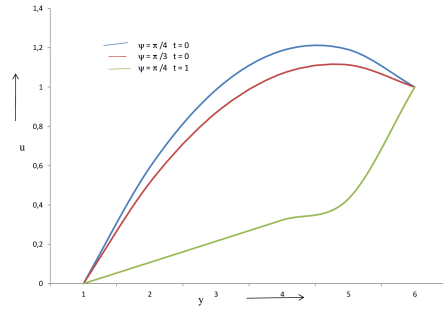


Figure 7: Velocity profile versus y for different of  $\psi = \pi/4$  and t when  $Ha=2, \lambda = 0.1, Re=2, S=5$   
 $R=2, K_1=1, Gr=10, Gc=1, Pe=0.71$  and  $K=1$

Figures 8-10 depicts that, Heat generation parameter S influences the temperature  $\theta$  whereas Radiation parameter R, Peclet number Pe, frequency  $\omega$  and t time suppresses.

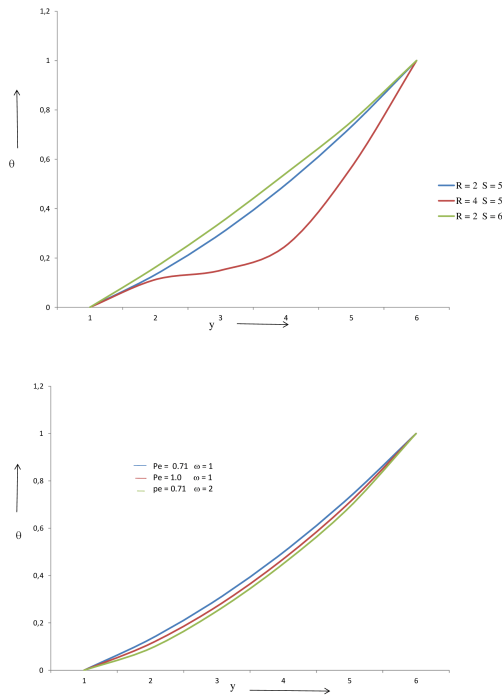


Figure 9: Temperature profile versus y for different values of Pe and  $\omega$  when  $R=2$  and  $S=5$ .

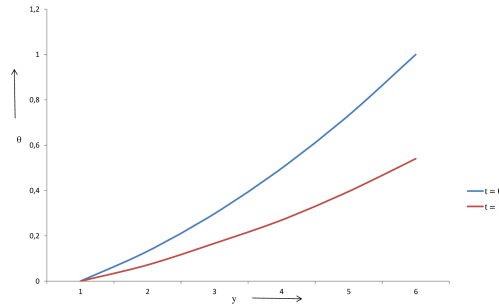


Figure 10: Temperature profile versus y for different of t when  $R=2$ ,  $S=5$ ,  $Pe=0.71$  and  $\omega=1$

Figures 11-12 shows that, Concentration  $\phi$  reduces for higher values of Schimdt number  $Sc$ , chemical reaction parameter  $Kr$  and frequency  $\omega$ .

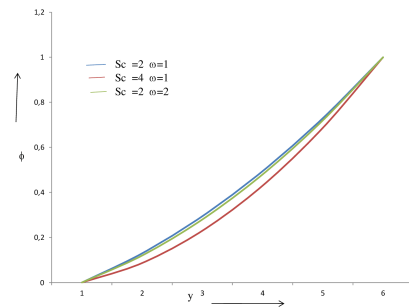


Figure 11: Concentration profile versus y for different of  $Sc$  and  $\omega$  when  $Kr=1, t=0$

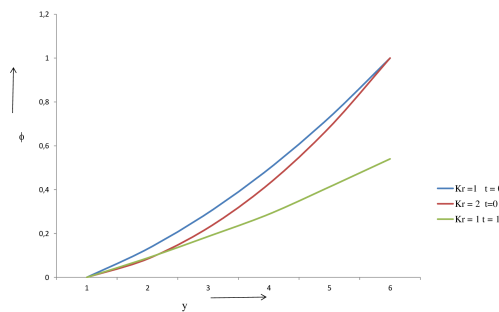


Figure 12: Concentration profile versus y for different of  $kr$  and  $t$  when  $Sc=2, \omega=1$

From Figure 13 , we draw conclusion that velocity decreases with increase in Casson parameter.

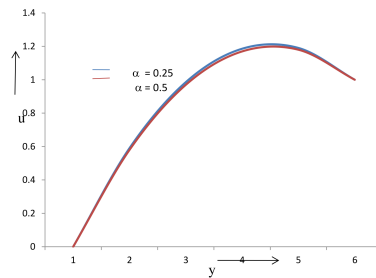


Figure 13 : Velocity profile versus y for different of Casson parameter w  
 $\lambda=0.1, Re=2, K=1, R=2, S=5, Pe=0.71, \omega=1, t=0, Kr=1, Gr=10, Gc=1$  and  $\psi =$

### Nomenclature:

$B_0$  Strength of the magnetic field

$C_p$  specific heat at constant pressure

$Gr$  Grashoff number

$Gc$  Solute Grashoff number

$g$  acceleration due to gravity

$Ha$  Hartmann Number

$K$  permeability parameter

$K^*$  permeability of porous medium

$P$  dimensionless pressure

$P^*$  fluid pressure

$Pe$  Peclet number

$Q_0$  heat generation/absorption constant

$q_r$  radiative heat flux in  $y^*$

$S$  heat source parameter

$T^*$  fluid temperature

$T_0$  Temperature of plate at  $y^*=0$

$T_w$  Temperature of plate at  $y^*=d$

$C_0$  Concentration of mass of plate at  $y^*=0$

$$\beta^* = 1 + \frac{1}{\alpha_1} \text{ Casson parameter}$$

$\alpha$ mean radiation absorption coefficient	
$R$ radiative parameter in the $y^*$ direction	$\beta$ coefficient of the thermal expansion
$R_e$ Reynolds number	$\theta$ dimensionless temperature
$\beta$ Coefficient of the thermal expansion	$\kappa$ thermal conductivity
$t$ dimensionless time	$\nu$ kinematic viscosity
$t^*$ time	$\sigma_e$ Electrical conductivity
$u$ dimensionless velocity along x-axis	$\rho$ fluid density
$u^*$ fluid velocity in $x^*$ direction	$\lambda$ wave length
$y$ dimensionless coordinate axis	$\psi$ ( $0 \leq \psi \leq \pi$ ) angle between velocity and strength of the magnetic field
$\beta^{**}$ Coefficient of mass concentration	

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